

towards Lax formulation of integrable hierarchies of topological type

Johan van de Leur

With:

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- ▶ Dubrovin-Zhang on Principal Hierarchies

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- ▶ Hirota biliner/quadratic equation
- ▶ Wave function and Sato-Wilson equation

KdV hierarchy, Lax equation:

$$\frac{\partial L}{\partial q_n} = \frac{a_n}{\sqrt{\hbar}} \left[\left(L^{n+\frac{1}{2}} \right)_+, L \right], \quad n = 0, 1, 2, \dots,$$

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$$\begin{aligned} L^{\frac{1}{2}} = & \sqrt{\hbar} \partial + \frac{u}{\sqrt{\hbar}} \partial^{-1} - \frac{u'}{2\sqrt{\hbar}} \partial^{-2} \\ & + \frac{\hbar u'' - 2u^2}{4\hbar^{3/2}} \partial^{-3} - \frac{\hbar u^{(3)} - 12uu'}{8\hbar^{3/2}} \partial^{-4} + \dots \end{aligned}$$

Lax equations are compatibility equations of the wave function ψ

$$\begin{aligned} L\psi(x, q_i, \lambda) &= \lambda^2 \psi(x, q_i, \lambda), \\ \frac{\partial \psi(x, q_i, \lambda)}{\partial q_n} &= \frac{a_n}{\sqrt{\hbar}} \left(L^{n+\frac{1}{2}} \right)_+ (\psi(x, q_i, \lambda)). \end{aligned}$$

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ψ satisfies the **Hirota bilinear/quadratic equation**:

$$\text{Res}_\lambda \lambda^{2p} \psi(x, q_i, \lambda) \psi(x', q'_i, -\lambda) = 0, \quad p = 0, 1, 2, \dots$$

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$$\Gamma_-(\lambda) = \exp \left(-\sqrt{\hbar} \sum_{n=0}^{\infty} b_n \lambda^{-2n-1} \frac{\partial}{\partial q_n} \right),$$

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- ▶ **Sato-Wilson equation:**

$$\frac{\partial P(\sqrt{h}\partial)}{\partial q_n} P(\sqrt{h}\partial)^{-1} = -\frac{a_n}{\sqrt{h}} \left(P(\sqrt{h}\partial) (\sqrt{h}\partial)^{2n+1} P(\sqrt{h}\partial)^{-1} \right)_-.$$

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Note:

$$\frac{a_n}{\sqrt{\hbar}} \text{Res}_\partial L^{n+\frac{1}{2}} = \frac{\partial^2 \log \tau}{\partial q_0 \partial q_n}$$

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 and $\tau(q) = \prod_{\gamma}^N \tau_{\gamma}(q_{\gamma 0}, q_{\gamma 1}, \dots)$
 then for every $1 \leq \alpha \leq N$:

$$\text{Res}_{\lambda_\alpha} \lambda_\alpha^{2p} \Gamma_\alpha(\lambda_\alpha) \tau(q) \otimes \Gamma_\alpha(-\lambda_\alpha) \tau(q') = 0$$

where

$$\Gamma_{\alpha,-}(\lambda_\alpha) = \exp \left(-\sqrt{\hbar} \sum_{n=0}^{\infty} b_n \lambda_\alpha^{-2n-1} \frac{\partial}{\partial q_{\alpha,n}} \right),$$

$$\Gamma_{\alpha,+}(\lambda_\alpha) = \exp \left(\frac{1}{\sqrt{\hbar}} \sum_{n=0}^{\infty} a_n q_{\alpha,n} \lambda_\alpha^{2n+1} \right)$$

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$$\hat{r} = \sum_{\alpha, \beta=1}^N (r_\ell)_{\alpha\beta} \left(\sum_{i=0}^{\infty} q_{\alpha, i} \frac{\partial}{\partial q_{\beta, \ell+i}} + \frac{\hbar}{2} \sum_{i+j=\ell-1} (-1)^{i+1} \frac{\partial^2}{\partial q_{\alpha, i} \partial q_{\beta, j}} \right),$$

$$\hat{s} = \sum_{\alpha, \beta=1}^N (s_\ell)_{\alpha\beta} \left(\sum_{i=0}^{\infty} q_{\alpha, i+\ell} \frac{\partial}{\partial q_{\beta, i}} - \frac{1}{2\hbar} \sum_{i+j=\ell-1} (-1)^{\ell+i} q_{\alpha, i} q_{\beta, j} \right).$$

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Let this $\hat{G}$ act on the Hirota bilinear identity:

$$\text{Res}_{\lambda_\alpha} \lambda_\alpha^{2p} (\hat{G} \Gamma_{\alpha,+}(\lambda_\alpha) \hat{G}^{-1}) (\hat{G} \Gamma_{\alpha,-}(\lambda_\alpha) \hat{G}^{-1}) \hat{G} \tau(q) \\ (\hat{G}' \Gamma_{\alpha,+}(-\lambda_\alpha) \hat{G}'^{-1}) (\hat{G} \Gamma_{\alpha,-}(-\lambda_\alpha) \hat{g}'^{-1}) \hat{G}' \tau(q') d\lambda_\alpha = 0.$$

Next rearrange, then up to $O(\epsilon^2)$ one finds:

$$\operatorname{Res}_\lambda \lambda^{2p} \left(G_{\alpha,+}(\lambda) \Gamma_{\alpha,+}(\lambda) G_{\alpha,-}(\lambda) \Gamma_{\alpha,-}(\lambda) \hat{G} \tau(q) \right. \\ \left. G'_{\alpha,+}(-\lambda) \Gamma_{\alpha,+}(-\lambda) G'_{\alpha,-}(-\lambda) \Gamma_{\alpha,-}(-\lambda) \hat{G}' \tau(q') \right) d\lambda = 0.$$

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where for $G = R$:

$$R_{\alpha,+}(\lambda) := \exp \left(\epsilon (r_\ell)_{\alpha\alpha} d_\ell \lambda^{2\ell} \right) \\ \times \exp \left(\frac{\epsilon (-1)^{\ell-1}}{\sqrt{\hbar}} \sum_{\beta} (r_\ell)_{\alpha\beta} \left(\sum_{n \geq \ell} a_n \lambda^{2n+1} q_{\beta,n-\ell} \right) \right) ; \\ R_{\alpha,-}(\lambda) := \exp \left(\epsilon \sqrt{\hbar} \sum_{\beta} (r_\ell)_{\alpha\beta} \sum_{n=0}^{\ell-1} (-1)^{n+1} a_n \lambda^{2n+1} \frac{\partial}{\partial q_{\beta,\ell-1-n}} \right) \\ \times \exp \left(\epsilon \sqrt{\hbar} \sum_{\beta, n \geq 0} (r_\ell)_{\alpha\beta} b_n \lambda^{-2n-1} \frac{\partial}{\partial q_{\beta,n+\ell}} \right).$$

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- ▶ divide this **new Hirota bilinear equation** by $\hat{G}\tau(q, x)\hat{G}'\tau(q', x')$
- ▶ and write the Hirota bilinear equation as:

$$\text{Res}_{\lambda} \lambda^{2p} V_{\alpha, G}(x, q, \lambda) e^{\frac{x\lambda}{\sqrt{h}}} V_{\alpha, G}(x', q', -\lambda) e^{\frac{-x'\lambda}{\sqrt{h}}} d\lambda = 0$$

where $V_{\alpha, G}(x, q, \lambda) =$

$$G_{\alpha, -}(\lambda)(P_{\alpha, G}(\lambda)) \frac{G_{\alpha, -}(\lambda)(\hat{G}\tau(x, q))}{\hat{G}\tau(x, q)} G_{\alpha, +}(\lambda) \Gamma_{\alpha, +}(\lambda) e^{\epsilon g_{\alpha, G}(x, \lambda)}$$

and

$$P_{\alpha, G}(\lambda) = \frac{\Gamma_{\alpha, -}(\lambda)(\hat{G}\tau(x, q))}{\hat{G}\tau(x, q)}$$

Write

$$V_{\alpha,G}(x, q, \sqrt{\hbar}\partial) = \left(P_{\alpha,G}(\sqrt{\hbar}\partial) + \epsilon Q_{\alpha,G}(x, q, \sqrt{\hbar}\partial) \right) \Gamma_{\alpha,+}(\sqrt{\hbar}\partial)$$

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Define

$$L_{\alpha,G} := P_{\alpha,G} \hbar \partial^2 P_{\alpha,G}^{-1} = L_{\alpha} + (L_{\alpha,G})_- = L_{\alpha} + \epsilon [L_{\alpha}, Q_{\alpha,G} P_{\alpha,G}^{-1}]_-$$

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$$L_{\alpha} = \hbar \partial^2 + 2u_{\alpha}(x, q, \epsilon)$$

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Deformed Sato-Wilson equation:

$$\begin{aligned} & \frac{\partial P_{\alpha,G}}{\partial q_{\beta,n}} P_{\alpha,G}^{-1} + \delta_{\alpha\beta} \frac{a_n}{\sqrt{\hbar}} \left(L_{\alpha,G}^{n+\frac{1}{2}} \right)_- \\ & + \epsilon \left(\frac{\partial Q_{\alpha,G} P_{\alpha,G}^{-1}}{\partial q_{\beta,n}} - \delta_{\alpha\beta} \frac{a_n}{\sqrt{\hbar}} \left[\left(L_{\alpha}^{n+\frac{1}{2}} \right)_+, Q_{\alpha,G} P_{\alpha,G}^{-1} \right] \right)_- = 0, \end{aligned}$$

$$\begin{aligned}
Q_{\alpha,R} P_{\alpha,R}^{-1} &= (r_\ell)_{\alpha\alpha} d_\ell L_{\alpha,R}^\ell + \frac{(-1)^{\ell-1}}{\sqrt{\hbar}} \sum_{\beta} (r_\ell)_{\alpha\beta} a_\ell x L_{\alpha,R}^{\ell+\frac{1}{2}} \\
&\quad - (r_\ell)_{\alpha\alpha} \sum_{n=0}^{\ell-1} (-1)^{n+1} a_n a_{\ell-1-n} \left(L_{\alpha,R}^{\ell-n-\frac{1}{2}} \right)_- L_{\alpha,R}^{n+\frac{1}{2}} \\
&\quad - (r_\ell)_{\alpha\alpha} \sum_{n \geq 0} b_n a_{n+\ell} \left(L_{\alpha,R}^{n+\ell+\frac{1}{2}} \right)_- L_{\alpha,R}^{-n-\frac{1}{2}} \\
&\quad + \frac{(-1)^{\ell-1}}{\sqrt{\hbar}} \sum_{\beta} (r_\ell)_{\alpha\beta} \sum_{n \geq \ell} a_n q_{\beta,n-\ell} L_{\alpha,R}^{n+\frac{1}{2}} \\
&\quad + \sqrt{\hbar} \sum_{\beta} (r_\ell)_{\alpha\beta} \sum_{n=0}^{\ell-1} (-1)^{n+1} a_n \frac{\partial \log \tau}{\partial q_{\beta,\ell-1-n}} L_{\alpha,R}^{n+\frac{1}{2}} \\
&\quad + \sqrt{\hbar} \sum_{\beta} (r_\ell)_{\alpha\beta} \sum_{n \geq \ell} b_n \frac{\partial \log \tau}{\partial q_{\beta,\ell+n}} L_{\alpha,R}^{-n-\frac{1}{2}}
\end{aligned}$$

In the work of Buryak, Posthuma and Shadrin:

$$\hat{G}\Omega_{\alpha,p;\beta,q} := \hbar \frac{\partial^2 \log \hat{G}\tau}{\partial q_{\alpha,p} \partial q_{\beta,q}}$$

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Now

$$\hat{G}\Omega_{\alpha,0;\beta,p} = -\text{Res}_{\partial} \frac{\partial P_{\alpha,G}}{\partial q_{\beta,q}} P_{\alpha,G}^{-1}$$

We show: The Res_{∂} of the deformed Sato-Wilson equations equal the deformation formula's of Buryak, Posthuma and Shadrin.