

The Gould-Hopper Polynomials in the Novikov-Veselov equation

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Outline

- 1 Introduction
- 2 Root Dynamics of σ -flows
- 3 Gould-Hopper Polynomials
- 4 Lax Pair of σ -flows
- 5 Smooth Rational Solutions of NV Equation

The Novikov-Veselov [1984] equation is defined by (U and t is **real**) [Athorne, Dubrovsky, Matveev, Nimmo, Salle, Y.Ohta]

$$\begin{aligned} U_t &= \partial_z^3 U + \bar{\partial}_z^3 U - 3\partial_z(VU) - 3\bar{\partial}_z(\bar{V}U), \\ \bar{\partial}_z V &= \partial_z U. \end{aligned} \quad (1)$$

When $z = \bar{z} = x$, we get the KdV equation ($U = V = \bar{V}$)

$$U_t = 2U_{xxx} - 12UU_x$$

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The Novikov-Veselov equation can be represented as the form of Manakov's triad

$$H_t = [A, H] - BH,$$

where H is the two-dimension Schrodinger operator

$$H = \partial_z \bar{\partial}_z + U$$

and

$$A = \partial_z^3 - V\partial_z + \bar{\partial}_z^3 - \bar{V}\bar{\partial}_z, \quad B = V_z + \bar{V}_{\bar{z}}.$$

It is equivalent to the linear representation

$$H\psi = 0, \quad \partial_t \psi = A\psi. \quad (2)$$

Motives:

- Integrable deformation of Schrodinger Operator [Athorne, Matveev, Nimmo, Y.Ohta]
- D-bar dressing method [Boiti, Dubrovsky, Leon, Pempinelli, Tsai]
- Two dimensional generalization of KdV equation [S.P.Novikov, A.Veselov, L.V. Bogdanov, P.G. Grinevich]
- Two-component BKP equation [R. Hirota, I.Krichever, Takasaki, Si-Qi Liu, C.Z. Wu, Y.J. Zhang]
- The Tzitzeica equation [E.Ferapontov., A.E. Mironov]
- Integrable deformation of minimal Lagrangian tori in CP^2 [A.E. Mironov]
- Integrable deformation of Dirichlet-to-Neumann map (Electrical Impedance Tomography, [M.Lassas, J. Mueller, A.Stahel])

Let $H\psi = H\omega = 0$. Then via the Moutard transformation[1878]

$$\begin{aligned} U(z, \bar{z}) &\longrightarrow \hat{U}(z, \bar{z}) = U(z, \bar{z}) + 2\partial\bar{\partial}\ln[i\int(\psi\partial\omega \\ &\quad - \omega\partial\psi)dz - (\psi\bar{\partial}\omega - \omega\bar{\partial}\psi)d\bar{z}], \\ \psi &\longrightarrow \theta = \frac{i}{\omega}\int(\psi\partial\omega - \omega\partial\psi)dz - (\psi\bar{\partial}\omega - \omega\bar{\partial}\psi)d\bar{z} \end{aligned}$$

one can construct a new Schrodinger operator

$$\hat{H} = \partial_z\bar{\partial}_z + \hat{U}$$

and $\hat{\psi} = \frac{1}{\theta}$ such that

$$\hat{H}\hat{\psi} = 0.$$

The **extended Moutard transformation** was established such that $\hat{U}(t, z, \bar{z})$ and $\hat{V}(t, z, \bar{z})$ defined by [Matveev, Salle, Athorne, Nimmo, H.C.Hu, S.Y.Lou and Q.P.Liu, 1991-2003]

$$\hat{U}(t, z, \bar{z}) = U(t, z, \bar{z}) + 2\partial\bar{\partial} \ln iW,$$

where

$$W = \int (\psi\partial\omega - \omega\partial\psi)dz - (\psi\bar{\partial}\omega - \omega\bar{\partial}\psi)d\bar{z} \quad (3)$$

$$+ [\psi\partial^3\omega - \omega\partial^3\psi + \omega\bar{\partial}^3 - \psi\bar{\partial}^3\omega \\ + 2(\partial^2\psi\partial\omega - \partial\psi\partial^2\omega) - 2(\bar{\partial}^2\psi\bar{\partial}\omega - \bar{\partial}\psi\bar{\partial}^2\omega) \quad (4)$$

$$+ 3V(\psi\partial\omega - \omega\partial\psi) - 3\bar{V}(\psi\bar{\partial}\omega - \omega\bar{\partial}\psi)]dt, \\ \hat{V} = V + 2\partial\bar{\partial} \ln iW, \quad (5)$$

will also satisfy the Novikov-Veselov equation.

In particular, we can use $U = V = 0$ as the seed solution. Then $H = \partial\bar{\partial}$. Let us consider the holomorphic functions $P(z, t)$:

$$\frac{\partial P}{\partial t} = \frac{\partial^3 P}{\partial z^3}. \quad (6)$$

Then we have the following

Theorem [Taimanov and Tsarev, 2008]

Let $\mathcal{P}_1(t, z)$ and $\mathcal{P}_2(t, z)$ be **polynomial functions** of z and satisfy (6). One defines $\omega_1 = \mathcal{P}_1 + \bar{\mathcal{P}}_1$ and $\omega_2 = \mathcal{P}_2 + \bar{\mathcal{P}}_2$. Then

$$U(t, z, \bar{z}) = 2\partial\bar{\partial} \ln iW,$$

where

$$\begin{aligned} W &= \mathcal{P}_1\bar{\mathcal{P}}_1 - \mathcal{P}_2\bar{\mathcal{P}}_2 + \int [(\mathcal{P}'_1\mathcal{P}_2 - \mathcal{P}_1\mathcal{P}'_2)dz \\ &\quad + (\bar{\mathcal{P}}_1\bar{\mathcal{P}}_2' - \bar{\mathcal{P}}_1'\bar{\mathcal{P}}_2)d\bar{z}] \\ &\quad + \int [\mathcal{P}_1'''\mathcal{P}_2 - \mathcal{P}_1\mathcal{P}_2''' + 2(\mathcal{P}'_1\mathcal{P}_2'' - \mathcal{P}_1''\mathcal{P}_2') + \bar{\mathcal{P}}_1\bar{\mathcal{P}}_2''' \\ &\quad - \bar{\mathcal{P}}_1'''\bar{\mathcal{P}}_2 + 2(\bar{\mathcal{P}}_1''\bar{\mathcal{P}}_2' - \bar{\mathcal{P}}_1'\bar{\mathcal{P}}_2''')]dt, \\ V &= 2\partial\bar{\partial} \ln iW, \end{aligned}$$

is a solution of Novikov-Veselov equation, which is **rational in**
 z, \bar{z}, t .

Taimanov and Tsarev considered the polynimial-type solutions for (6)

$$P_N(t, z) = z^N + \sigma_1 z^{N-1} + \sigma_2 z^{N-2} + \cdots + \sigma_{N-1} z + \sigma_N.$$

Then the flow (6) generates the σ -flow:

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$$P_N(t, z) = z^N + \sigma_1 z^{N-1} + \sigma_2 z^{N-2} + \cdots + \sigma_{N-1} z + \sigma_N.$$

Then the flow (6) generates the σ -flow:

$$\dot{\sigma}_k = (N - K + 3)(N - k + 2)(N - k + 1)\sigma_{k-3}, k = 1, 2, 3 \cdots N.$$

It can be seen that σ_1, σ_2 are conserved quantities. Indeed, $\sigma_1, \sigma_2, \cdots, \sigma_N$ are the elementary symmetric polynomials in the roots q_1, q_2, \cdots, q_N of $P(z)$:

$$\begin{aligned} \sigma_1(\vec{q}) &= - \sum_{i=1}^N q_i, & \sigma_2(\vec{q}) &= \sum_{i < j} q_i q_j, \\ \sigma_3(\vec{q}) &= - \sum_{i < j < k} q_i q_j q_k, \cdots, & \sigma_N(\vec{q}) &= (-1)^N q_1 q_2 \cdots q_N. \end{aligned} \quad (7)$$

The integrable (even linear) evolution of $\vec{\sigma} = (\sigma_1, \sigma_2, \dots, \sigma_N)$ induces a dynamical system on the symmetric product $S^N C$ of the complex roots . We call such a **dynamical system on $S^N C$ a σ -system**.

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Remark:

From (7), we see that given two solutions $\mathcal{P}_1(t, z)$ and $\mathcal{P}_2(t, z)$, by a substitution of **$e^{i\lambda_1} \mathcal{P}_1(t, z)$ and $e^{i\lambda_2} \mathcal{P}_2(t, z)$** , where λ_1 and λ_2 are **real-valued constants**, into (7) we obtain a solution of the Novikov-Veselov solution. Therefore, to each pair of holomorphic solutions of (6), we can get an **$(S^1 \times S^1)$** -family of solutions to the Novikov-Veselov equation.

Let's write $P_N(t, z)$ as

$$P_N(t, z) = (z - q_1(t))(z - q_2(t)) \cdots (z - q_N(t)).$$

Then from the equation (6), one gets the **root dynamics**

$$\dot{q}_j = -6 \sum_{\substack{m < n, \\ j \neq m, n}}^N \frac{1}{(q_j - q_m)(q_j - q_n)}. \quad (8)$$

For example, when $N=3$, we have

$$\begin{aligned} \dot{q}_1 &= -6 \frac{1}{(q_1 - q_2)(q_1 - q_3)} \\ \dot{q}_2 &= -6 \frac{1}{(q_2 - q_1)(q_2 - q_3)} \\ \dot{q}_3 &= -6 \frac{1}{(q_3 - q_1)(q_3 - q_2)} \end{aligned}$$

For $N=4$, one has

$$\begin{aligned}\dot{q}_1 &= -6\left[\frac{1}{(q_1 - q_2)(q_1 - q_3)} + \frac{1}{(q_1 - q_3)(q_1 - q_4)} + \frac{1}{(q_1 - q_2)(q_1 - q_4)}\right] \\ \dot{q}_2 &= -6\left[\frac{1}{(q_2 - q_1)(q_2 - q_3)} + \frac{1}{(q_2 - q_3)(q_2 - q_4)} + \frac{1}{(q_2 - q_1)(q_2 - q_4)}\right] \\ \dot{q}_3 &= -6\left[\frac{1}{(q_3 - q_1)(q_3 - q_2)} + \frac{1}{(q_3 - q_1)(q_3 - q_4)} + \frac{1}{(q_3 - q_2)(q_3 - q_4)}\right] \\ \dot{q}_4 &= -6\left[\frac{1}{(q_4 - q_2)(q_4 - q_3)} + \frac{1}{(q_4 - q_1)(q_4 - q_2)} + \frac{1}{(q_4 - q_1)(q_1 - q_3)}\right]\end{aligned}$$

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We notice that since σ_1 and σ_2 are conserved quantities, one knows that

$$\sum_{i=1}^N q_i, \quad \sum_{i=1}^N q_i^2$$

are conserved densities of (8)

The goal is to investigate the properties of the root dynamics (8):

- Initial Value Problem
- Lax Pair
- Asymptotic behavior

The generating function of the Gould-Hopper polynomials $P_N(t, z)$ is

$$e^{\lambda z + \lambda^3 t} = \sum_{N=0}^{\infty} P_N(t, z) \frac{\lambda^N}{N!}.$$

Indeed, the Gould-Hopper polynomials $P_N(t, z)$ has the operator representation [1962]

$$P_N(t, z) = e^{t\partial_z^3} z^N = [1 + t\partial_z^3 + \frac{t^2\partial_z^6}{2!} + \frac{t^3\partial_z^9}{3!} + \frac{t^4\partial_z^{12}}{4!} + \cdots] z^N.$$

We remark that in general the Gould-Hopper polynomials are defined by $P_N^{(m)}(t, z) = e^{t\partial_z^m} z^N$. Here we take $m = 3$.

One notices that the Gould-Hopper polynomials $P_N(t, z)$ are characterized by (6) and $P_N(0, z) = z^N$. For example,

$$\begin{aligned} P_0 &= 1, & P_1 &= z, & P_2 &= z^2, & P_3 &= z^3 + 6t, & P_4 &= z^4 + 24tz, \\ P_5 &= z^5 + 60tz^2, & P_6 &= z^6 + 120tz^3 + 360t^2, \\ P_7 &= z^7 + 210tz^4 + 2520z, & P_8 &= z^8 + 336tz^5 + 10080t^2z^2, \\ P_9 &= z^9 + 504tz^6 + 30240t^2z^3 + 60480t^3, \\ P_{10} &= z^{10} + 720tz^7 + 75600t^2z^4 + 604800t^3z \end{aligned}$$

Actually, we have

$$\begin{aligned} P_N(t, z) &= N! \sum_{k=0}^{[N/3]} \frac{t^k z^{N-3k}}{k!(N-3k)!} \\ \frac{dP_N(t, z)}{dz} &= NP_{N-1}(t, z) \end{aligned} \tag{9}$$

From the operation calculus, one has

$$(z + 3t\partial_z^2)P_N(t, z) = P_{N+1}(t, z).$$

Hence we yield the recursive relation

$$P_{N+1}(t, z) = zP_N(t, z) + 3t(N-1)(N-2)P_{N-3}(t, z). \quad (10)$$

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We notice that if we consider the equation (6) with the initial data of analytical function

$$P(0, z) = \sum_{N=0}^{\infty} \alpha_N z^N,$$

then the formal solution is

$$P(t, z) = e^{t\partial_z^3} \sum_{N=0}^{\infty} \alpha_N z^N = \sum_{N=0}^{\infty} \alpha_N P_N(t, z). \quad (11)$$

Remark:

The **successive operations** of the operator $(z + 3t\partial_z^2)$ on the solution (11) can help us construct more solutions of (6). For example, if $P(0, z) = \sin z$, then we have,

$$\begin{aligned} e^{t\partial_z^3} \sin z &= e^{t\partial_z^3} \sum_{N=0}^{\infty} \frac{(-1)^N}{(2N+1)!} z^{2N+1} = \sum_{N=0}^{\infty} \frac{(-1)^N}{(2N+1)!} P_{2N+1}(t, z) \\ &= \sin(z - t). \end{aligned}$$

The last equation uses the fact $e^{t\partial_z^3} e^{iz} = e^{iz+t}$. Hence

$$(z + 3t\partial_z^2)^N \sin(z - t), \quad N = 0, 1, 2, 3, 4, \dots$$

are also solutions of (6).

- Initial Value Problem : The root dynamics of σ -flow can be solved by

$$\begin{aligned} H_N(t, z) &= (z - q_1(t))(z - q_2(t)) \cdots (z - q_N(t)) \\ &= P_N(t, z) + C_1 P_{N-1}(t, z) + \cdots + C_N P_0(t, z), \end{aligned}$$

where the constants $C_1, C_2, \dots, C_{N-1}, C_N$ are determined by the initial values of $q_1(0), q_2(0), \dots, q_N(0)$, that is,

$$\begin{aligned} C_1 &= - \sum_{i=1}^N q_i(0), & C_2 &= \sum_{i < j} q_i(0) q_j(0), \\ C_3 &= - \sum_{i < j < k} q_i(0) q_j(0) q_k(0), & \cdots, \\ C_N &= (-1)^N q_1(0) q_2(0) \cdots q_N(0). \end{aligned}$$

It can be seen that the solutions $q_1(t), q_2(t), \dots, q_N(t)$ can be obtained **algebraically**.

- Lax pair:

Firstly, we study the root dynamics of the Gould-Hopper polynomials, which correspond to the initial values

$$q_1(0) = q_2(0) = \cdots = q_N(0) = 0.$$

Let's define the $N \times N$ matrix by

$$X(t) = \begin{cases} a_{i,i+1} = 1, & \text{if } i = 1, 2, 3, \dots, N; \\ a_{i,i-2} = -3t(i-1)(i-2), & \text{if } i = 3, 4, \dots, N-1; \\ 0, & \text{otherwise.} \end{cases} \quad (12)$$

Then from the recursive relation (10), one knows that

$$P_N(t, z) = \det(zI_N - X(t)).$$

For example, when $N = 3$,

$$X(t) = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ -6t & 0 & 0 \end{pmatrix};$$

N=4,

$$X(t) = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ -6t & 0 & 0 & 1 \\ 0 & -18t & 0 & 0 \end{pmatrix};$$

N=5,

$$X(t) = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ -6t & 0 & 0 & 1 & 0 \\ 0 & -18t & 0 & 0 & 1 \\ 0 & 0 & -36t & 0 & 0 \end{pmatrix}.$$

We can write $X(t)$ as

$$X(t) = R(t)QR^{-1}(t),$$

where $Q = \text{diag}(q_1(t), q_2(t), \dots, q_N(t))$ and

$$R(t) = \begin{pmatrix} P_0(q_1, t) & P_0(q_2, t) & P_0(q_3, t) & \cdots & P_0(q_N, t) \\ P_1(q_1, t) & P_1(q_2, t) & P_1(q_3, t) & \cdots & P_1(q_N, t) \\ P_2(q_1, t) & P_2(q_2, t) & P_2(q_3, t) & \cdots & P_2(q_N, t) \\ \vdots & & & & \\ P_N(q_1, t) & P_N(q_2, t) & P_N(q_3, t) & \cdots & P_N(q_N, t) \end{pmatrix}. \quad (13)$$

For instance, when $N = 3$,

$$R(t) = \begin{pmatrix} 1 & 1 & 1 \\ q_1 & q_2 & q_3 \\ q_1^2 & q_2^2 & q_3^2 \end{pmatrix};$$

$N=4$,

$$R(t) = \begin{pmatrix} 1 & 1 & 1 & 1 \\ q_1 & q_2 & q_3 & q_4 \\ q_1^2 & q_2^2 & q_3^2 & q_4^2 \\ q_1^3 + 6t & q_2^3 + 6t & q_3^3 + 6t & q_4^3 + 6t \end{pmatrix};$$

$N=5$,

$$R(t) = \begin{pmatrix} 1 & 1 & 1 & 1 & 1 \\ q_1 & q_2 & q_3 & q_4 & q_5 \\ q_1^2 & q_2^2 & q_3^2 & q_4^2 & q_5^2 \\ q_1^3 + 6t & q_2^3 + 6t & q_3^3 + 6t & q_4^3 + 6t & q_5^3 + 6t \\ q_1^4 + 24t & q_2^4 + 24t & q_3^4 + 24t & q_4^4 + 24t & q_5^4 + 24t \end{pmatrix}.$$

From the initial value problem, we notice that the **polynomials t^n** can be replaced by the **elementary symmetric polynomials of the roots q_1, q_2, \dots, q_N** . Hence one has $R(\vec{q})$. It can be seen that

$$\dot{X}(t) = RLR^{-1},$$

where

$$L = \dot{Q} + [M, Q], \quad M = R^{-1}\dot{R}.$$

For example, when $N=3$,

$$L(t) = \begin{pmatrix} \dot{q}_1 & \dot{q}_2 \frac{q_2 - q_3}{q_3 - q_1} & \dot{q}_3 \frac{q_3 - q_2}{q_2 - q_1} \\ \dot{q}_1 \frac{q_1 - q_3}{q_3 - q_2} & \dot{q}_2 & \dot{q}_3 \frac{q_3 - q_1}{q_1 - q_2} \\ \dot{q}_1 \frac{q_1 - q_2}{q_2 - q_3} & \dot{q}_2 \frac{q_2 - q_1}{q_1 - q_3} & \dot{q}_3 \end{pmatrix};$$

$N=4$,

$$L(t) = \begin{pmatrix} \dot{q}_1 & -\frac{\dot{q}_2(q_2 - q_3)(q_2 - q_4) + 6}{(q_1 - q_3)(q_1 - q_4)} & -\frac{\dot{q}_3(q_3 - q_2)(q_3 - q_4) + 6}{(q_1 - q_2)(q_1 - q_4)} & -\frac{\dot{q}_4(q_4 - q_2)(q_4 - q_3) + 6}{(q_1 - q_2)(q_1 - q_3)} \\ -\frac{\dot{q}_1(q_1 - q_3)(q_1 - q_4) + 6}{(q_2 - q_3)(q_2 - q_4)} & \dot{q}_2 & -\frac{\dot{q}_3(q_3 - q_1)(q_3 - q_4) + 6}{(q_2 - q_4)(q_2 - q_1)} & -\frac{\dot{q}_4(q_4 - q_1)(q_4 - q_3) + 6}{(q_2 - q_3)(q_2 - q_1)} \\ -\frac{\dot{q}_1(q_1 - q_2)(q_1 - q_4) + 6}{(q_2 - q_3)(q_4 - q_3)} & -\frac{\dot{q}_2(q_2 - q_4)(q_2 - q_1) + 6}{(q_4 - q_3)(q_1 - q_3)} & \dot{q}_3 & -\frac{\dot{q}_4(q_4 - q_1)(q_4 - q_2) + 6}{(q_2 - q_3)(q_1 - q_3)} \\ -\frac{\dot{q}_1(q_1 - q_2)(q_1 - q_3) + 6}{(q_4 - q_3)(q_4 - q_2)} & -\frac{\dot{q}_2(q_2 - q_3)(q_2 - q_1) + 6}{(q_4 - q_3)(q_4 - q_1)} & -\frac{\dot{q}_3(q_3 - q_1)(q_3 - q_2) + 6}{(q_4 - q_1)(q_4 - q_2)} & \dot{q}_4 \end{pmatrix};$$

Since

$$\frac{dX(t)}{dt} = \begin{cases} a_{i,i-2} = -3(i-1)(i-2), & \text{if } i = 3, 4, \dots, N-1; \\ 0, & \text{otherwise,} \end{cases}$$

we know $\frac{dX(t)}{dt}$ is a nilpotent matrix and hence L is a nilpotent one, too. So

$$\text{tr}(L^r) = \text{tr}\left[\frac{dX(t)}{dt}\right]^r = 0, \quad r = 1, 2, 3, \dots, \dots$$

Actually, a simple calculation yields

$$L^{[\frac{N}{2}]+1} = 0, \quad N \geq 3.$$

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Actually, a simple calculation yields

$$L^{[\frac{N}{2}]+1} = 0, \quad N \geq 3.$$

Now,

$$\frac{d^2 X(t)}{dt^2} = 0$$

will imply the Lax equation

$$\frac{dL(t)}{dt} = [L, M]. \tag{14}$$

For $N = 3, 4, 5$, we see that, by Maple software, q_i satisfies the following Goldfish model [Calogero, 2001], a limiting case of the Ruijsenaars-Schneider system:

$$\ddot{q}_i = 2 \sum_{j \neq i} \frac{\dot{q}_i \dot{q}_j}{q_i - q_j}.$$

The reason is that the Gould-Hopper polynomials P_N , $N = 3, 4, 5$ are **linear in t -variable**.

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The reason is that the Gould-Hopper polynomials P_N , $N = 3, 4, 5$ are **linear in t -variable**. For $N = 6$, we have from the diagonal terms of the Lax equation (14)

$$\ddot{q}_i = 2 \sum_{j \neq i}^6 \frac{\dot{q}_i \dot{q}_j}{q_i - q_j} + \frac{\sum_{j=1}^6 (\text{some quadratic terms of } \vec{q}) q_j + 720}{\prod_{i \neq j}^6 (q_i - q_j)}.$$

Remark:
For the Goldfish Model

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its initial value problem can be solved by the statement:

$z = q_i(t), i = 1, 2, \dots, N$ are the N roots of the equation

$$\sum_{i=1}^N \frac{\dot{q}_i(0)}{z - q_i(0)} = \frac{1}{t}.$$

It can be seen that it is a **polynomial in z with coefficients linear in t** . Then the special choices of initial datum can get the solutions of the root dynamics (8) for the cases **$N = 3, 4, 5$** .

Secondly, we consider the general case. Let's define the 2D Appell polynomials $\mathbb{G}_n(z, t)$ by means of the generating function [G.Bretti and P.Ricci, 2004]:

$$A(\lambda)e^{\lambda z + \lambda^3 t} = \sum_{n=0}^{\infty} \mathbb{G}_n(z, t) \frac{\lambda^n}{n!}, \quad (15)$$

where

$$A(\lambda) = \sum_0^N \frac{\Gamma_k}{k!} \lambda^k,$$

Γ'_k 's being constants and $\Gamma_0 = 1$. Then one has the following formula, noting that $\binom{N}{h} = \binom{N}{N-h}$,

$$\mathbb{G}_N = \sum_{h=0}^N \binom{N}{N-h} \Gamma_{N-h} P_h(z, t). \quad (16)$$

It's easy to see that the polynomials $\mathbb{G}_n(z, t)$ also satisfy the linear equation (6).

Then we have

$$\mathbf{\Gamma}_{N-h} = \frac{C_{N-h}}{\binom{N}{N-h}}. \quad (17)$$

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Now, it's suitable to introduce the coefficients of the Taylor expansion

$$\frac{A'(\lambda)}{A(\lambda)} = \sum_{n=0}^{\infty} \alpha_n \frac{\lambda^n}{n!}.$$

It can be seen that the coefficients α_n can be expressed by $\Gamma_0, \Gamma_1, \dots, \Gamma_{n+1}$ (or the initial values (17)). For examples,

$$\begin{aligned} \alpha_0 &= \Gamma_1, & \alpha_1 &= \Gamma_2 - \Gamma_1^2, & \alpha_2 &= \Gamma_3 - 3\Gamma_1\Gamma_2 + 2\Gamma_1^3, \\ \alpha_3 &= \Gamma_4 + 12\Gamma_1^2\Gamma_2 - 4\Gamma_1\Gamma_3 - 3\Gamma_2^2 - 6\Gamma_1^4, \\ &\vdots \end{aligned}$$

The recurrence relation for the 2D Appell polynomial $\mathbb{G}_N(z, t)$ can be written as follows:

$$\begin{aligned}\mathbb{G}_0(z, t) &= 1 \\ \mathbb{G}_N(z, t) &= (z + \alpha_0)\mathbb{G}_{N-1}(z, t) + 3t(N-1)(N-2)\mathbb{G}_{N-3}(z, t) \\ &\quad + \sum_{k=0}^{N-2} \binom{N-1}{k} \alpha_{N-k-1} \mathbb{G}_k(z, t).\end{aligned}\tag{18}$$

A simple calculation can yield

$$\begin{aligned}\mathbb{G}_0(z, t) &= 1, \quad \mathbb{G}_1(z, t) = z + \alpha_0, \quad \mathbb{G}_2(z, t) = (z + \alpha_0)^2 + \alpha_1 \\ \mathbb{G}_3(z, t) &= (z + \alpha_0)^3 + 3\alpha_1(z + \alpha_0) + \alpha_2 + 6t \\ \mathbb{G}_4(z, t) &= (z + \alpha_0)^4 + 6\alpha_1(z + \alpha_0)^2 + (4\alpha_2 + 24t)(z + \alpha_0) + \alpha_3 \\ &\quad + 3\alpha_1^2.\end{aligned}$$

When $A(\lambda) = 1$, this recursive relation becomes (10). Hence the relation (18) is a generalization of (10) for arbitrary initial data.

The matrix corresponding to (12) can be constructed as follows:

$$X(t) = \begin{cases} a_{i,i+1} = 1, & \text{if } i = 1, 2, 3, \dots, N-1; \\ a_{i,i} = -\alpha_0, & \text{if } i = 1, 2, 3, \dots, N; \\ a_{i,i-1} = -\binom{i-1}{i-2} \alpha_1, & \text{if } i = 2, 3, 4, \dots, N; \\ a_{i,i-2} = -\binom{i-1}{i-3} (6t + \alpha_2), & \text{if } i = 3, 4, 5, \dots, N; \\ a_{i,i-3} = -\binom{i-1}{i-4} \alpha_3, & \text{if } i = 4, 5, 6, \dots, N; \\ \vdots & \\ a_{i,i-k} = -\binom{i-1}{i-k-1} \alpha_k, & \text{if } i = k+1, k+2, k+3, \dots, N; \\ \vdots & \\ a_{N,1} = -\alpha_{N-1} \end{cases}$$

Similarly, one has

$$\mathbb{G}_N(z, t) = \det(zI_N - X(t)).$$

For instance, when $N = 5$, we get

$$X(t) = \begin{pmatrix} -\alpha_0 & 1 & 0 & 0 & 0 \\ -\alpha_1 & -\alpha_0 & 1 & 0 & 0 \\ -(6t + \alpha_2) & -2\alpha_1 & -\alpha_0 & 1 & 0 \\ -\alpha_3 & -(18t + 3\alpha_2) & -3\alpha_1 & -\alpha_0 & 1 \\ -\alpha_4 & -4\alpha_3 & -(36t + 6\alpha_2) & -4\alpha_1 & -\alpha_0 \end{pmatrix}.$$

Also, one can write $X(t)$ as

$$X(t) = R(t)QR^{-1}(t),$$

where $Q = \text{diag}(q_1(t), q_2(t), \dots, q_N(t))$ and $R(t)$ is defined as (13) with $P_m(q_i, t)$ being replaced by $\mathbb{G}_m(q_i, t)$. Then one follows the previous procedures and finally can get the Lax equation (14) for general case. Therefore the root dynamics (8) is Lax-integrable.

We notice here that for $N = 3, 4, 5$ the root dynamics of \mathbb{G}_N also satisfies the Gold-fish model .

- Asymptotic behavior

It is known that the Gould-Hopper polynomial $P_N(t, z)$ has the scaling property:

$$P_N(t, z) = t^{\frac{N}{3}} \hat{P}_N\left(\frac{z}{t^{1/3}}\right), \quad (19)$$

where $\hat{P}_N(\eta)$ is the so-called **generalized Hermite polynomial (or Appell polynomial)** in $\eta = \frac{z}{t^{1/3}}$. For example,

$$\begin{aligned} P_8(t, z) &= z^8 + 336tz^5 + 10080t^2 = t^{\frac{8}{3}}[\eta^8 + 336\eta^5 + 10080\eta^2] \\ &= t^{\frac{8}{3}}\hat{P}_8(\eta). \end{aligned}$$

Then the k -th zero $\lambda_N^{(k)}$ (**constant**) of $\hat{P}_N(\eta)$ determines the dynamics of the root q_k , i.e.,

$$q_k(t) = \lambda_N^{(k)} t^{1/3}.$$

Since $\hat{P}_N(\xi \lambda_N^{(k)}) = 0$, $\xi^3 = 1$, one knows that the roots q_k are located on the circles in the plane with time dependent radius or fixed at the origin.

Since $\hat{P}_N(\xi \lambda_N^{(k)}) = 0, \xi^3 = 1$, one knows that the roots q_k are located on the **circles in the plane with time dependent radius or fixed at the origin**. Finally, from the Initial value Problem, we know that when $t \rightarrow \infty$ and $z \rightarrow \infty$ such that $|z|^3/t \rightarrow \text{constant}$, $P_N(t, z)$ plays the dominant role. Hence one yields

$$q_k(t) \rightarrow \lambda_N^{(k)} t^{1/3}.$$

Consequently, the roots asymptotically will follow diagonal lines.

In this section, we establish **smooth rational solutions** for all time by the extended Moutard transformation (7) and Gould-Hopper polynomials.

Example 1

Let $\mathcal{P}_1 = z^2 + z + 1$ and $\mathcal{P}_2 = -iz^2 - 2iz$. Then a simple calculation can yield the imaginary part of W in (7)

$$M(x, y, t) = (x^2 + y^2)^2 + \frac{8}{3}x^3 + 4xy^2 + 4x^2 + 4x + 4t + 100.$$

We can see that

$$M(x, y, 0) \approx 4x + 100 \quad \text{near} \quad (0, 0)$$

and

$$M(x, y, 0) \approx (x^2 + y^2)^2 \quad \text{near} \quad r = \sqrt{x^2 + y^2} = \infty.$$

It can be verified that $M(x, y, 0)$ is positive for all \mathbb{R}^2 . Also, **$M(x, y, t)$ is positive for all \mathbb{R}^2 at any fixed time t** . Then the solution U of the Novikov-Veselov equation (1) is

where

$$\begin{aligned} M_1 = & -12(294 + 600x^2 + 588y^2 + 8x^3 + 888x + 12t + 3x^4 \\ & - 6x^2y^2 - 2x^3y^2 + 24x^2t - 3y^4x + 24y^2t + 36xt - 3y^4 + x^5) \end{aligned}$$

and

$$\begin{aligned} M_2 = & (3x^4 + 6x^2y^2 + 3y^4 + 8x^3 + 12xy^2 + 12x^2 \\ & + 12x + 12t + 300)^2 \end{aligned}$$

At fixed time t , one knows U decays like $\frac{1}{r^3}$ for $r \rightarrow \infty$. Also, U tends asymptotically to zero at the rate $\frac{1}{t}$ at any fixed point (x, y) when $t \rightarrow \infty$.

Example 2

Let $\mathcal{P}_1 = (z^3 + 6t) + 2iz$ and $\mathcal{P}_2 = -i(z^3 + 6t) + z$. Then a simple calculation can yield the imaginary part of W in (7)

$$\begin{aligned} f(x, y, t) &= (x^2 + y^2)^3 + 4x^3y + 8xy^3 + 2(x^2 + y^2) \\ &+ 6t(2x^3 - 6xy^2 - y) + 36t^2 + 6000. \end{aligned}$$

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Then we can see that

$$f(x, y, 0) \approx 2(x^2 + y^2) + 6000 \quad \text{near} \quad (0, 0)$$

and

$$f(x, y, 0) \approx (x^2 + y^2)^3 \quad \text{near} \quad r = \sqrt{x^2 + y^2} = \infty.$$

It can be verified that $f(x, y, 0)$ is positive for all \mathbb{R}^2 .

Also, letting $f(x, y, t)$ be equal to zero, one has

$$t = \frac{(1/2)xy^2 + (1/12)y - (1/6)x^3 \pm (1/12)\sqrt{24x^2y^4 - 20xy^3 - 36x^4y^2 - 7y^2 - 20x^3y - 4y^6 - 8x^2 - 23996}}{-8x^2 - 23996}.$$

A simple calculation shows that the equation inside the square root is negative for all \mathbb{R}^2 . Hence $f(x, y, t)$ is positive for all \mathbb{R}^2 at any fixed time t . Then a solution U of the Novikov-Veselov equation (1) is

$$U = \frac{F_1}{F_2},$$

where

$$\begin{aligned}
F_1 = & -\frac{1}{2}[24x^7y + 16x^6 + 24x^5y^3 + (432ty + 216000 + 48y^2)x^4 \\
& + (-24y^5 - 144t + 48y)x^3 + (432000y^2 + 864ty^3 - 96y^4)x^2 \\
& + (-24y^7 - 48y^3 + 432ty^2 + (432000 + 1728t^2)y)x + 48000 \\
& + 432y^5t - 32y^6 + 216000y^4 + 252t^2]
\end{aligned}$$

and

$$\begin{aligned}
F_2 = & (x^6 + 3x^4y^2 + 3x^2y^4 + y^6 + 4x^3y + 8xy^3 + 2x^2 + 2y^2 \\
& + 12tx^3 - 36txy^2 - 6ty + 36t^2 + 6000)^2
\end{aligned}$$

At fixed time t , one knows U decays like $\frac{1}{r^4}$ for $r \rightarrow \infty$. Also, U tends asymptotically to zero at the rate $\frac{1}{t^2}$ at any fixed point (x, y) when $t \rightarrow \infty$.

Example 3

Let $\mathcal{P}_1 = (z^5 + 60tz^2) + 2iz$ and $\mathcal{P}_2 = -i(z^5 + 60tz^2) + z$. Then the imaginary part of W in (7) is

$$\begin{aligned} & h(x, y, t) \\ = & \left((x^2 + y^2)^5 - \frac{10}{3} x^5 y + \frac{20}{3} x^3 y^3 \right. \\ & + (x^2 + y^2) (2 + 12 x^3 y - 12 x y^3) \\ & + 120 t (x^2 + y^2) (x^5 - 3 y^4 x - 2 x^3 y^2) + 3600 t^2 (x^2 + y^2)^2 \\ & \left. + 20 t (11 y^3 + 3 x^2 y) + 120 t^2 + 1000 \right). \end{aligned}$$

Similarly, $h(x, y, t)$ is positive for all \mathbb{R}^2 at fixed time t . Then the solution U is

$$U = \frac{H_1}{H_2},$$

where

$$\begin{aligned}
& H_1(x, y, t) \\
= & -\frac{1}{2}[360xy^{13} + (360x^3 + 14400t)y^{11} - 144y^{10} + (14400x^2t - 1800x^5)y^9 \\
& + (-720x^2 + 900000 + 108000t^2)y^8 + (237600tx^4 - 777600t^2x - 3600x^7)y^7 \\
& + (3840x^4 + (3600000 + 432000t^2)x^2 + 83520tx)y^6 + (165600tx^6 + 720x \\
& - 10368000t^3 - 1800x^9 - 345600t^2x^3)y^5 + (-1120x^6 + (5400000 \\
& + 648000t^2)x^4 + 12000tx^3 + (-129600000t - 15552000t^3)x - 810000t^2)y^4 \\
& + (360x^{11} - 50400x^8t + 950400x^5t^2 - 2160x^3 + 20736000t^3x^2 + (-216000t^2 \\
& - 1800000)x - 5760t)y^3 + (2380x^8 + (3600000 + 432000t^2)x^6 - 25200tx^5 \\
& + (-10368000t^3 - 86400000t)x^3 + 453600x^2t^2 + 62208000t^4 + 518400000t^2)y^2 \\
& + (360x^{13} + 21600x^{10}t + 518400x^7t^2 + 720x^5 + 31104000t^3x^4 + (1920000 \\
& + 230400t^2)x^3 + 17280x^2t + 1555200t^3 + 12960000t)y + 476x^{10} \\
& + (900000 + 108000t^2)x^8 + 25200x^7t + (43200000t + 5184000t^3)x^5 \\
& + 226800x^4t^2 + (62208000t^4 + 518400000t^2)x^2 + 72000 + 8640t^2)]
\end{aligned}$$

and

$$\begin{aligned}
& H_2(x, y, t) \\
&= (3x^{10} + 15x^8y^2 + 30x^6y^4 + 30y^6x^4 + 15y^8x^2 \\
&+ 3y^{10} + 26x^5y + 20x^3y^3 \\
&+ 6x^2 + 6y^2 - 36xy^5 + 360x^7t - 1800tx^3y^4 - 360tx^5y^2 - 1080txy^6 \\
&+ 10800x^4t^2 + 21600t^2x^2y^2 + 10800t^2y^4 + 660ty^3 + 180tx^2y \\
&+ 360t^2 + 3000)^2
\end{aligned}$$

In this case, at fixed time, U decays like $\frac{1}{r^6}$ for $r \rightarrow \infty$;
 moreover, we see that

$$U \rightarrow \frac{-240(x^2 + y^2)}{[30(x^2 + y^2)^2 + 1]^2} = \frac{-240z\bar{z}}{[30z^2\bar{z}^2 + 1]^2} \quad \text{as } t \rightarrow \infty \quad (20)$$

at fixed point (x, y) , which is a **stationary solution** of (1) for

$$V = \frac{3600\bar{z}^4z^2 - 120\bar{z}^2}{(30z^2\bar{z}^2 + 1)^2}.$$

We notice that if we define

$$u(z, \bar{z}) = \ln U \quad \text{and} \quad V(z, \bar{z}) = \frac{U_{zz}}{3U},$$

then the **stationary equation** of the Novikov-Veselov equation (1) will become the **Tzitzeica equation**[B.C. Konopelchenko and U. Pinkall., 1998]

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$$u_{z\bar{z}} = e^u + \epsilon e^{-2u}, \quad (21)$$

where ϵ is an arbitrary constant. It can be verified that (20) satisfies the **$\epsilon = 0$ case, i.e., the Liouville equation**, whose real solutions are given by

$$\ln \frac{-2\kappa \left| \frac{dS}{dz} \right|^2}{(1 + \kappa |S|^2)^2}, \quad (22)$$

where $S(z)$ is a **locally univalent meromorphic function in some domain** and κ is a constant. For the solution corresponding to (20), one knows that $S(z) = z^2$ and $\kappa = 30$.

For general case, let's choose

$$\mathcal{P}_1 = P_N(t, z) + 2iz \quad \text{and} \quad \mathcal{P}_2 = -iP_N(t, z) + z.$$

We can expect the imaginary part of W in (7) is positive for all \mathbb{R}^2 at any fixed time **if we choose an appropriate constant**. And the solution $U(x, y, t)$ at any fixed time **decays like $\frac{1}{r^{N+1}}$ for $N \geq 2$** . When $t \rightarrow \infty$ and then one obtains the solutions (22) of the Liouville equation or the ones of the Tzitzeica equation (21).

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Remark: For each potential $U(x, y, t)$ there exist **infinitely many** wave functions, which can be constructed by the **Pfaffian** of linear combinations of the Gould-Hopper polynomials

In summary, one investigates

- The Gould-Hopper polynomials
- The σ -flow and the Lax pair
- The Gold-Fish model
- The smooth rational solutions of NV equation
- The solution of the Liouville equation (or the Tzitzeica equation)

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Thanks for your attention.