

Frobenius Manifold for the Dispersionless Kadomtsev–Petviashvili Equation

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- I Frobenius manifolds and systems of quasilinear PDE
- II dKP equation and Vlasov equations
- III Frobenius manifold for dKP
- IV Principal hierarchy and solutions of dKP

Frobenius manifolds

Frobenius manifold (Dubrovin, ['92]):

- M : differentiable manifold, $\dim(M) = n$,
- η : flat, non-degenerate, (pseudo-)metric
- $\circ$: commutative, associative product
- e : unity vector field: $X \circ e = X, \forall X$.
- E : Euler vector field

Notation:

$$c(X, Y, Z) := \eta(X, Y \circ Z)$$

$$c^i_{jk} := \eta^{is} c_{sjk}, \longrightarrow (X \circ Y)^i = X^j c^i_{jk} Y^k$$

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Frobenius manifolds

In addition, we require:

(1) Compatibility with the metric (Frobenius algebra on $T_p M$):

$$\eta(X, Y \circ Z) = \eta(X \circ Y, Z)$$

(2) Compatibility with the connection

$$\nabla_W c(X, Y, Z) = \nabla_X c(W, Y, Z)$$

∇ : covariant derivative of η .

Remark

(1) + (2) $\implies \exists F \in C^\infty(M)$ s.t.

$$c_{ijk} = \nabla_i \nabla_j \nabla_k F$$

$F =$ *potential* $of the Frobenius manifold (\implies WDVV \dots)$

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Moreover:

(3) Flatness of unity vector field:

$$\nabla e = 0.$$

(4) Euler vector field satisfies

$$\nabla \nabla E = 0$$

$$[E, X \circ Y] - [E, X] \circ Y - X \circ [E, Y] = X \circ Y$$

$$E(\eta(X, Y)) - \eta([E, X], Y) - \eta(X, [E, Y]) = (2 - d) \eta(X, Y)$$

d = charge of the Frobenius manifold

Frobenius manifolds – further properties

i) $\exists$ a second metric: **intersection form** of the Frobenius manifold

$$g^{ij} = E^k \eta^{is} c_{ks}^j,$$

g is **flat** and $\eta + \alpha g$ is **flat** $\forall \alpha$ (compatible flat metrics)

ii) Special coordinate sets:

a) **Flat** coordinates: η is constant

b) **Canonical** coordinates: $c_{jk}^i = \delta_j^i \delta_k^i$ (*semisimple Frob. manifold*)

*In this talk: **no** flat, **no** canonical coordinates!*

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Systems of quasilinear PDE

Systems of quasilinear PDEs: $u^i = u^i(x, t)$, $i = 1, \dots, n$

$$u_t^i = V_j^i(u) u_x^j, \quad \lim_{|x| \rightarrow \infty} u^i(x, t) = 0.$$

Hamiltonian form: (Dubrovin, Novikov ['83])

$$\{\mathcal{H}_1, \mathcal{H}_2\} = \int \frac{\delta \mathcal{H}_1}{\delta u^i(x)} \left(\eta^{ij} \frac{\partial}{\partial x} - \eta^{is} \Gamma_{sk}^j u_x^k \right) \frac{\delta \mathcal{H}_2}{\delta u^j(x)} dx, \quad \mathcal{H}_\alpha = \int H_\alpha(u) dx$$

η flat pseudo-metric, Γ Christoffel symbols

Integrability: Generalized hodograph method (Tsarev, ['94]).

$$x \delta_j^i + t V_j^i(u) = W_j^i(u), \quad u_\tau^i = W_j^i(u) u_x^j, \quad \frac{\partial^2 u^i}{\partial t \partial \tau} = \frac{\partial^2 u^i}{\partial \tau \partial t}$$

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Principal hierarchy of a Frobenius manifold

$(M, \eta, \mathbf{c}, e, E)$ Frobenius manifold, X vector field on M

$$\eta \implies (\bar{V}_X)_j^i := \nabla_j X^i, \quad \mathbf{c} \implies (V_X)_j^i := c_{jk}^i X^k$$

Recursive relation. For $\alpha = 1, \dots, n$:

$$\nabla_j X_{\alpha,0}^i = 0$$

$$\nabla_j X_{\alpha,m+1}^i = c_{jk}^i X_{\alpha,m}^k, \quad m = 0, 1, 2, \dots$$

Principal hierarchy of a Frobenius manifold

$$\frac{\partial u^i}{\partial t_{\alpha,m}} = (V_{\alpha,m})_j^i \frac{\partial u^j}{\partial x}, \quad (V_{\alpha,m})_j^i = c_{jk}^i X_{\alpha,m}^k,$$

The flows of the principal hierarchy pairwise commute:

$$\frac{\partial^2 u^i}{\partial t_{\beta,m} \partial t_{\alpha,n}} = \frac{\partial^2 u^i}{\partial t_{\beta,m} \partial t_{\alpha,n}}, \quad \forall i, \alpha, \beta, n, m$$

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The flows of the principal hierarchy *pairwise commute*:

$$\frac{\partial^2 u^i}{\partial t_{\beta,m} \partial t_{\alpha,n}} = \frac{\partial^2 u^i}{\partial t_{\beta,m} \partial t_{\alpha,n}}, \quad \forall i, \alpha, \beta, n, m$$

dKP and Vlasov equations

Let $f(p, x, t, y)$ be a solution of the **integro - differential** equations (**Vlasov**, or **collisionless Boltzmann** equation):

$$f_y = p f_x - A_x^0 f_p, \quad (1)$$

$$f_t = (p^2 + A^0) f_x - (A_x^0 p + A_x^1) f_p, \quad (2)$$

where

$$A^k(x, y, t) = \int_{-\infty}^{+\infty} p^k f(p, x, y, t) dp, \quad k \in \mathbb{N}.$$

The flows (1), (2) commute: $(f_y)_t = (f_t)_y$.

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dKP and Vlasov equations

Moment equations:

$$A_y^k = A_x^{k+1} + kA^{k-1}A_x^0, \quad k \in \mathbb{N} \quad (\text{Benney chain})$$

$$A_t^k = A_x^{k+2} + A^0 A_x^k + (k+1)A^k A_x^0 + kA^{k-1}A_x^1, \quad k \in \mathbb{N}$$

Important fact: $A^0 = \int f \, dp$ satisfies the dKP equation:

$$A_{yy}^0 = \partial_x (A_t^0 - A^0 A_x^0)$$

Indeed:

$$A_y^0 = A_x^1, \quad A_y^1 = A_x^2 + A^0 A_x^0, \quad A_t^0 = A_x^2 + 2A^0 A_x^0.$$

Rearranging:

$$\begin{aligned} A_x^1 &= A_y^0, \\ A_y^1 &= A_t^0 - A^0 A_x^0, \end{aligned}$$

which is compatible.

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dKP and Vlasov equations – Hamiltonian structure

Poisson–Vlasov bracket (Marsden, Weinstein, [’82]):

$$\{\mathcal{G}, \mathcal{H}\}_{LP} := \iint f(x, p) \left\{ \frac{\delta \mathcal{G}[f]}{\delta f(x, p)}, \frac{\delta \mathcal{H}[f]}{\delta f(x, p)} \right\}_{x,p} dp dx,$$

$\{\cdot, \cdot\}_{x,p}$: canonical bracket, $\mathcal{G}, \mathcal{H}$ functionals of f .

Hamilton's equations:

$$\frac{\partial f}{\partial t} = \{f, \mathcal{H}\}_{LP}, \quad \Longleftrightarrow \quad \frac{\partial f}{\partial t} + \left\{ f, \frac{\delta \mathcal{H}}{\delta f} \right\}_{x,p} = 0,$$

Example

$$\mathcal{H} = \frac{1}{2} \iint (p^2 + A^0) f dp dx \quad \Longrightarrow \quad \frac{\partial f}{\partial y} = \left\{ f, \frac{p^2}{2} + A^0 \right\}_{x,p}$$

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dKP – Lax formulation

Formal series

$$\lambda := p + \sum_{k=0}^{\infty} \frac{A^k}{p^{k+1}}, \quad A^k = A^k(x + t_1, t_2, t_3, \dots).$$

The **flows** (dispersionless Lax equations):

$$\lambda_{t_n} = \{\Omega_n, \lambda\}_{x,p}, \quad \Omega_n = \frac{1}{n} (\lambda^n)_+, \quad k \in \mathbb{N},$$

Zero curvature equations (Zakharov-Shabat):

$$\partial_{t_n} \Omega_m - \partial_{t_m} \Omega_n = \{\Omega_m, \Omega_n\}_{x,p}$$

Remark (bi-Hamiltonian structure)

dKP admits **two** compatible structures of Dubrovin–Novikov type

$$\text{First metric:} \quad \eta^{kn}(A) = (k+n)A^{k+n-1}$$

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Following (Gibbons, ['81], Gibbons, Tsarev ['94,'96]): more analytical structure on λ . We consider:

$$\lambda(p) = p + \int_{-\infty}^{+\infty} \frac{f(q)}{p - q} dq, \quad p \in \mathbb{R},$$

Asymptotic expansion

$$\lambda \sim p + \sum_{k=0}^{\infty} \frac{A^k}{p^{k+1}}, \quad p \mapsto \infty$$

where

$$A^k = \int_{-\infty}^{+\infty} p^k f(p) dp, \quad k \in \mathbb{N},$$

are the moments of f .

Known examples of infinite dimensional Frobenius manifolds

[Krichever, Mineev-Weinstein, Wiegmann, Zabrodin, A. \['04\]](#)

Solutions of WDVV depending on infinitely many coordinates and associated to the dispersionless $2D$ Toda hierarchy. [Zabrodin \['09\]](#): dKP hierarchy

[Carlet, Dubrovin, Mertens \['10\]](#): Frobenius manifold for dispersionless $2D$ Toda hierarchy on the space of pairs of functions analytic inside and outside the unit circle respectively, and with prescribed singularities at zero and at infinity.

[Wu, Xu \['11\]](#): Frobenius manifold for the dispersionless two-component BKP hierarchy

Vlasov eqs as Hydrodynamic type systems

Following ([Gibbons, A.R. \['07\]](#)): we consider Vlasov equations as 'continuous indexed' hydrodynamic type systems:

$$f_t(p) = \int V\left(\frac{p}{q}\right) f_x(q) dq$$

Naive idea:

$$\{u^1, \dots, u^n\} \longrightarrow \{f(p), p \in \mathbb{R}\}$$

$$\{u^1(x), \dots, u^n(x)\} \longrightarrow \{f(p, x), p \in \mathbb{R}\}$$

$$\sum_i \longrightarrow \int dp$$

$$\frac{\partial}{\partial u^i} \longrightarrow \frac{\delta}{\delta f(p)}$$

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The Schwartz class

Analytical properties of f in the (x, p) -plane:

p -dependence: f in **Schwartz class** $\mathcal{S}(\mathbb{R})$ of $\mathcal{C}^\infty$ -functions s.t.

$$\sup_{p \in \mathbb{R}} |p^k \partial^m f(p)| < +\infty, \quad k, m \in \mathbb{N}$$

$$f \in \mathcal{S}(\mathbb{R}) \implies \lambda(p) = p + \int \frac{f(q)}{p-q} dq \in \mathcal{C}^\infty(\mathbb{R})$$

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Look for linear maps

$$V\left(\begin{smallmatrix} p \\ q \end{smallmatrix}\right) : \mathcal{S}(\mathbb{R}) \longrightarrow \mathcal{S}(\mathbb{R}),$$

thus thinking at $\mathcal{S}(\mathbb{R})$ as space of **tangent vectors**.

Space of **cotangent vectors**: dual space of $\mathcal{S}(\mathbb{R})$ through the L^2 -pairing

$$\langle \omega, \Phi \rangle = \int_{-\infty}^{+\infty} \omega(p) \Phi(p) dp, \quad \omega \in \mathcal{S}'(\mathbb{R}), \Phi \in \mathcal{S}(\mathbb{R}).$$

$\mathcal{S}'(\mathbb{R})$, space of **tempered distributions**.

Example ($V\left(\begin{smallmatrix} p \\ q \end{smallmatrix}\right) = p \delta(p - q) - f'(p)$)

$$\int V\left(\begin{smallmatrix} p \\ q \end{smallmatrix}\right) f_x(q) dq = p f_x(p) - A_x^0 f'(p) \in \mathcal{S}(\mathbb{R})$$

Later on: we will consider **proper** subspaces of $\mathcal{S}(\mathbb{R})$ and $\mathcal{S}'(\mathbb{R})$

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*Later on: we will consider **proper** subspaces of $\mathcal{S}(\mathbb{R})$ and $\mathcal{S}'(\mathbb{R})$*

The Schwartz class

Look for linear maps

$$V\left(\begin{smallmatrix} p \\ q \end{smallmatrix}\right) : \mathcal{S}(\mathbb{R}) \longrightarrow \mathcal{S}(\mathbb{R}),$$

thus thinking at $\mathcal{S}(\mathbb{R})$ as space of **tangent vectors**.

Space of **cotangent vectors**: dual space of $\mathcal{S}(\mathbb{R})$ through the L^2 -pairing

$$\langle \omega, \Phi \rangle = \int_{-\infty}^{+\infty} \omega(p) \Phi(p) dp, \quad \omega \in \mathcal{S}'(\mathbb{R}), \Phi \in \mathcal{S}(\mathbb{R}).$$

$\mathcal{S}'(\mathbb{R})$, space of **tempered distributions**.

Example ($V\left(\begin{smallmatrix} p \\ q \end{smallmatrix}\right) = p \delta(p - q) - f'(p)$)

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Later on: we will consider **proper** subspaces of $\mathcal{S}(\mathbb{R})$ and $\mathcal{S}'(\mathbb{R})$

First contravariant metric: $\eta^{(p,q)}$

Proposition (Gibbons, A.R., ['07])

The LP-bracket can be written as a continuous index bracket of hydrodynamic type

$$\{\mathcal{G}, \mathcal{H}\} = \iiint \frac{\delta \mathcal{G}[f]}{\delta f(x, p)} \left(\eta^{(p,q)} \frac{\partial}{\partial x} + \int \Gamma^{(p,q)}_r f_x(r) dr \right) \frac{\delta \mathcal{H}[f]}{\delta f(x, p)} dp dq dx,$$

Continuous form of the first metric

$$\eta^{(p,q)}[f] = -f'(p)\delta(p-q) \implies \eta^{nm}[A] = (n+m)A^{n+m-1}$$

Remark

η is not invertible on $S'(\mathbb{R})$. We consider $\mathcal{O}_M(\mathbb{R}) \subset S'(\mathbb{R})$, the space of $C^\infty(\mathbb{R})$ functions which, together with all their derivatives, grow at infinity slower than some polynomial.

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First covariant metric: $\eta_{(p,q)}$

Denote $\mathcal{V} \subset \mathcal{S}(\mathbb{R})$ the **linear** space of functions of the form

$$X^{(p)} = h(p)f'(p), \quad h \in \mathcal{O}_M(\mathbb{R})$$

On $\mathcal{V}$, the **inverse metric** of η :

$$\eta_{(p,q)} = -\frac{1}{f'(p)}\delta(p-q),$$

is well defined. Action on vectors $X_1, X_2 \in \mathcal{V}$:

$$\eta(X_1, X_2) = \iint X_1^{(p)} \eta_{(p,q)} X_2^{(q)} dp dq = - \int \frac{X_1^{(p)} X_2^{(p)}}{f'(p)} dp.$$

*The metric is **flat**: the (continuous) Riemann tensor is **zero***

Christoffel symbols

Christoffel symbols

$$\Gamma\left(\begin{smallmatrix} p \\ q r \end{smallmatrix}\right) = \frac{1}{f'(q)} \delta'(q - p) \delta(q - r).$$

Covariant derivative, acting on $X \in \mathcal{V}$ as

$$\nabla_{(q)} X^{(p)} = \frac{\delta X^{(p)}}{\delta f(q)} + \int \Gamma\left(\begin{smallmatrix} p \\ q r \end{smallmatrix}\right) X^{(r)} dr = \frac{\delta X^{(p)}}{\delta f(q)} + \delta'(q - p) X^{(p)},$$

and on $h \in \mathcal{O}_M(\mathbb{R})$ as

$$\nabla_{(p)} h_{(q)} = \frac{\delta h_{(q)}}{\delta f(p)} - \int \Gamma\left(\begin{smallmatrix} r \\ p q \end{smallmatrix}\right) h_{(r)} dr = \frac{\delta h_{(q)}}{\delta f(p)} - \delta(p - q) \frac{\partial h_{(p)}}{\partial p}.$$

The two actions can be proved to be **consistent**, and can be extended to higher order multilinear maps.

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The two actions can be proved to be **consistent**, and can be extended to higher order multilinear maps.

Structure constants of the algebra

Look for a bilinear map $\circ : \mathcal{V} \times \mathcal{V} \longrightarrow \mathcal{V}$, of the form

$$(X \circ Y)^{(p)} = \iint X^{(q)} c\left(\begin{smallmatrix} p \\ qs \end{smallmatrix}\right) Y^{(s)} dq ds, \quad X, Y \in \mathcal{V}$$

Consider the quantities

$$c\left(\begin{smallmatrix} p \\ qr \end{smallmatrix}\right) = \frac{\delta(p-r)}{p-q} - \frac{f'(p)}{f'(q)} \frac{\delta(q-r)}{p-q} + \frac{\delta(p-q)}{q-r} - \frac{\lambda'(p)}{f'(p)} \delta(q-r) \delta(q-p),$$

where

$$\lambda'(p) = 1 + \int_{-\infty}^{+\infty} \frac{f'(q)}{p-q} dq.$$

The space $\mathcal{V}$ is **closed** under the product $\circ$. Moreover, we have:

Structure constants of the algebra – properties

- Symmetry

$$c\left(\begin{smallmatrix} p \\ q \ r \end{smallmatrix}\right) = c\left(\begin{smallmatrix} p \\ r \ q \end{smallmatrix}\right).$$

- Associativity

$$\int c\left(\begin{smallmatrix} p \\ q \ s \end{smallmatrix}\right) c\left(\begin{smallmatrix} s \\ l \ r \end{smallmatrix}\right) ds = \int c\left(\begin{smallmatrix} p \\ l \ s \end{smallmatrix}\right) c\left(\begin{smallmatrix} s \\ q \ r \end{smallmatrix}\right) ds.$$

- Compatibility with the metric

$$\int \eta_{(p \ q)} c\left(\begin{smallmatrix} q \\ r \ s \end{smallmatrix}\right) dq = \int \eta_{(r \ q)} c\left(\begin{smallmatrix} q \\ p \ s \end{smallmatrix}\right) dq.$$

- Compatibility with the connection

$$\nabla_{(l)} c\left(\begin{smallmatrix} p \\ q \ r \end{smallmatrix}\right) = \nabla_{(q)} c\left(\begin{smallmatrix} p \\ l \ r \end{smallmatrix}\right).$$

The unity vector: e

Proposition

The element $e^{(p)} = -f'(p) \in \mathcal{V}$ is the **unity** vector field.

Indeed:

$$(X \circ e)^{(p)} = \iint X^{(q)} c\left(\begin{smallmatrix} p \\ q r \end{smallmatrix}\right) e^{(r)} dq dr = \int X^{(q)} \delta(p - q) dq = X^{(p)}$$

Moreover,

$$\nabla_{(q)} e^{(p)} = \frac{\delta e^{(p)}}{\delta f(q)} + \int \Gamma\left(\begin{smallmatrix} p \\ q s \end{smallmatrix}\right) e^{(s)} ds = -\delta'(p - q) - \delta'(q - p) = 0,$$

$\implies e$ is **flat**.

Proposition

The element

$$E^{(p)} = f(p) - p f'(p) \in \mathcal{S}(\mathbb{R}) \setminus \mathcal{V}$$

satisfies

$$\nabla_p \nabla_q E^{(r)} = 0,$$

$$\begin{aligned} & \int E^{(s)} \frac{\delta c_{(qr)}^{(p)}}{\delta f(s)} ds - \int c_{(qr)}^{(s)} \frac{\delta E^{(p)}}{\delta f(s)} ds + \int \frac{\delta E^{(s)}}{\delta f(q)} c_{(sr)}^{(p)} ds + \int \frac{\delta E^{(s)}}{\delta f(r)} c_{(qs)}^{(p)} ds = c_{(qr)}^{(p)} \\ & \int E^{(r)} \frac{\delta \eta_{(pq)}}{\delta f(r)} dr + \int \eta_{(rq)} \frac{\delta E^{(r)}}{\delta f(p)} dr + \int \eta_{(pq)} \frac{\delta E^{(r)}}{\delta f(q)} dr = 3 \eta_{(pq)}, \end{aligned}$$

and it is therefore the **Euler vector field** of the Frobenius manifold. Moreover, the Frobenius manifold has charge **$d = -1$** .

Potential of the Frobenius manifold

Theorem

The functional (*logarithmic energy with external field*)

$$F = \frac{1}{2} \iint \log |p - q| f(p) f(q) dp dq + \frac{1}{2} \int p^2 f(p) dp.$$

satisfies the condition

$$\nabla_{(p)} \nabla_{(q)} \nabla_{(r)} F = c_{(pqr)}$$

and is therefore the potential of the Frobenius manifold.

*Random Matrix Theory: equilibrium measure for the large N limit of the partition function for the Gaussian Unitary Ensemble.
(Wiegmann, Zabrodin ['00], Elbau, Felder ['05], ...)*

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PDE associated with the product c : dKP hierarchy

Example $(X_2^{(p)} = -p f'(p) \in \mathcal{V})$

$$\begin{aligned} V_{X_2} \left(\frac{p}{q} \right) &:= \int c \left(\frac{p}{q}, r \right) X_2^{(r)} dr \\ &= -\frac{p f'(p)}{p - q} + \frac{q f'(p)}{p - q} - \delta(p - q) \left(\int \frac{r f'(r)}{p - r} dr - p \lambda'(p) \right) \\ &= -f'(p) - \delta(p - q) \left(\int \frac{(r - p) f'(r)}{p - r} dr - p \right) \\ &= p \delta(p - q) - f'(p), \end{aligned}$$

Example $(X_3^{(p)} = -(p^2 + 2A^0) f'(p) \in \mathcal{V})$

$$V_{X_3} \left(\frac{p}{q} \right) = (p^2 + A^0) \delta(p - q) - (p + q) f'(p)$$

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Principal hierarchy – Flat vector fields

The functionals

$$\mathcal{H}_{h,0} = \int H_{h,0} dx, \quad H_{h,0} = \int h(f) dp,$$

are **Casimirs** of the Lie–Poisson bracket. Consider the vector fields

$$X_{h,0}^{(p)} = \int \frac{\delta H_{h,0}}{\delta f(q)} \eta^{(q,p)} dq = -h'(f(p))f'(p) \in \mathcal{V}, \quad (3)$$

Proposition

The vector fields (3) are flat:

$$\nabla_{(q)} X_{h,0}^{(p)} = 0.$$

For $h(f) = f$ we get the unity $e^{(p)} = -f'(p)$.

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Principal Hierarchy – Primary flows

Primary flows of the principal hierarchy:

$$\partial_{t_{h,0}} f(p) = \int V_{h,0} \left(\frac{p}{q} \right) f_x(q) dq, \quad V_{h,0} \left(\frac{p}{q} \right) := \int c \left(\frac{p}{qr} \right) X_{h,0}^{(r)} \quad (4)$$

Proposition

The flows (4) are Hamiltonian of the form

$$\partial_{t_{h,0}} f(p) = \left\{ f(p), \int H_{h,1} dx \right\}_{LP},$$

with the Poisson–Vlasov bracket, and the Hamiltonian

$$H_{h,1} = \int h(f(p)) \lambda(p) dp.$$

Theorem (Principal Hierarchy for dKP)

The vector fields

$$X_{h,n}^{(p)} := -\frac{\delta H_{h,n}}{\delta f(p)} f'(p), \quad H_{h,n} = \frac{1}{n!} \int h(f(p)) \lambda(p)^n dp,$$

$n \in \mathbb{N}$, satisfy the recurrence relations

$$\nabla_{(q)} X_{h,n+1}^{(p)} = \int c\left(\begin{smallmatrix} p \\ qr \end{smallmatrix}\right) X_{h,n}^{(r)},$$

of the principal hierarchy. The corresponding commuting flows

$$\partial_{t_{h,n}} f(p) = \int V_{h,n}\left(\begin{smallmatrix} p \\ q \end{smallmatrix}\right) f_x(q) dq, \quad V_{h,n}\left(\begin{smallmatrix} p \\ q \end{smallmatrix}\right) := \int c\left(\begin{smallmatrix} p \\ qr \end{smallmatrix}\right) X_{h,n}^{(r)},$$

are Hamiltonian

$$\partial_{t_{h,n}} f(p) = \{f, \mathcal{H}_{h,n+1}\}_{LP}, \quad \mathcal{H}_{h,n+1} = \int H_{h,n+1} dx.$$

The generic conserved quantity

Proposition

Let $k(\mu, \nu)$, be a function of two variables, s.t. the integral

$$\mathcal{K} = \iint k(f(p, x), \lambda(p, x)) dp dx,$$

converges. Then, $\mathcal{K}$ is a conserved quantity for the dKP hierarchy.

Direct result: no dependence on the Frobenius manifold!!

$$\text{Principal hierarchy} \implies h(f, \lambda) = \frac{1}{n!} h(f) \lambda^n$$

$$\text{Classical Hamiltonians} \implies h(f, \lambda) = \frac{1}{n!} f \lambda^n + C(f)$$

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Hodograph formula for dKP

Analogy with finite dimensional case:

$$x V_{X_1} \left(\frac{p}{q} \right) + y V_{X_2} \left(\frac{p}{q} \right) + t V_{X_3} \left(\frac{p}{q} \right) - V_{X_K} \left(\frac{p}{q} \right) = 0,$$

$$\text{with } X_1 = -f'(p), \quad X_2 = -pf'(p), \quad X_3 = -(p^2 + 2A^0)f'(p)$$

$$K = \int k(f(p), \lambda(p)) dp \implies X_K = -\frac{\delta K}{\delta f(r)} f'(r)$$

$$\int c \left(\frac{p}{qr} \right) \left(x + y r + t (r^2 + 2A^0) - \frac{\delta K}{\delta f(r)} \right) f'(r) dr = 0,$$

Nonlinear integral equation (also in Manakov, Santini [06]):

$$x + y p + t (p^2 + 2A^0) = \partial_1 k(f(p), \lambda(p)) - \int \frac{\partial_2 k(f(q), \lambda(q))}{p - q} dq$$

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Hodograph formula – Solutions of dKP

Hodograph formula = stationary point of a dKP conserved quantity

Indeed: the functional

$$\mathcal{H}_K = \int \left(x A^0 + y A^1 + t \left(A^2 + (A^0)^2 \right) - K \right) dx.$$

is a **constant of motion** for the y and t flows of the hierarchy, and

$$\text{Hodograph formula} \iff \frac{\delta \mathcal{H}_K}{\delta f(p, x)} = 0.$$

Moreover:

$$f \text{ satisfies Hodograph} \implies A^0 = \int f dp \text{ is a solution of dKP}$$

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